# **Physics**

# NCERT Exemplar Problems

### Chapter 8

# Electromagnetic Waves

### Answers

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$\mathbf{o}$ . $\mathbf{I}$	(C)

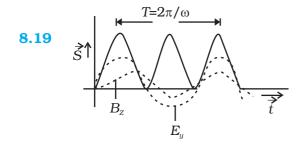
- **8.2** (b)
- **8.3** (b)
- **8.4** (d)
- **8.5** (d)
- **8.6** (c)
- **8.7** (c)
- **8.8** (a), (d)
- **8.9** (a), (b), (c)
- **8.10** (b), (d)

- **8.14** As electromagnetic waves are plane polarised, so the receiving antenna should be parallel to electric/magnetic part of the wave.
- **8.15** Frequency of the microwave matches the resonant frequency of water molecules.

**8.16** 
$$i_C = i_D = \frac{dq}{dt} = -2\pi q_0 v \sin 2\pi v t$$
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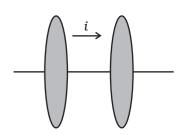
8.17 On decreasing the frequency, reactance  $X_c = \frac{1}{\omega C}$  will increase which will lead to decrease in conduction current. In this case  $i_D = i_C$ ; hence displacement current will decrease.

8.18 
$$I_{av} = \frac{1}{2}c\frac{B_0^2}{\mu_0} = \frac{1}{2} \times \frac{3 \times 10^8 \times (12 \times 10^{-8})^2}{1.26 \times 10^{-6}} = 1.71 W/m^2.$$



**8.20** EM waves exert radiation pressure. Tails of comets are due to solar solar radiation.

8.21 
$$B = \frac{\mu_0 2I_D}{4\pi r} = \frac{\mu_0 1}{4\pi r} = \frac{\mu_0}{2\pi r} \varepsilon_0 \frac{d\phi_E}{dt}$$
$$= \frac{\mu_0 \varepsilon_0}{2\pi r} \frac{d}{dt} (E\pi r^2)$$
$$= \frac{\mu_0 \varepsilon_0 r}{2} \frac{dE}{dt}.$$



**8.22** (a)  $\lambda_1 \to \text{Microwave}, \quad \lambda_2 \to \text{UV}$   $\lambda_3 \to \text{X rays}, \quad \lambda_4 \to \text{Infrared}$ 

(b) 
$$\lambda_3 < \lambda_2 < \lambda_4 < \lambda_1$$

Microwave - Radar
 UV - LASIK eye surgery
 X-ray - Bone fracture identification (bone scanning)
 Infrared - Optical communication.

8.23 
$$S_{av} = c^{2} \varepsilon_{0} | \mathbf{E}_{0} \times \mathbf{B}_{0} | \frac{1}{T} \int_{0}^{T} \cos^{2}(kx - \omega t) dt \text{ as } \mathbf{S} = c^{2} \varepsilon_{0} (\mathbf{E} \times \mathbf{B})$$

$$= c^{2} \varepsilon_{0} E_{0} B_{0} \frac{1}{T} \times \frac{T}{2}$$

$$= c^{2} \varepsilon_{0} E_{0} \left( \frac{E_{0}}{c} \right) \times \frac{1}{2} \left( as c = \frac{E_{0}}{B_{0}} \right)$$

$$= \frac{1}{2} \varepsilon_{0} E_{0}^{2} c$$

$$= \frac{E_{0}^{2}}{2\mu_{0}c} as \left( c = \frac{1}{\sqrt{\mu_{0} \varepsilon_{0}}} \right)$$

$$\mathbf{8.24} \qquad i_D = C \frac{dV}{dt}$$

$$1 \times 10^{-3} = 2 \times 10^{-6} \frac{dV}{dt}$$

$$\frac{dV}{dt} = \frac{1}{2} \times 10^3 = 5 \times 10V / s$$

Hence, applying a varying potential difference of  $5 \times 10^2$  V/s would produce a displacement current of desired value.

**8.25** Pressure

$$P = \frac{Force}{Area} = \frac{F}{A} = \frac{1}{A} \frac{\Delta p}{\Delta t} \quad (F = \frac{\Delta p}{\Delta t} = \text{rate of change of momentum})$$

$$= \frac{1}{A} \cdot \frac{U}{\Delta tc} \quad (\Delta pc = \Delta U = \text{energy imparted by wave in time} \Delta t)$$

$$= \frac{I}{c} \left( \text{intensity } I = \frac{U}{A\Delta t} \right)$$

8.26 Intensity is reduced to one fourth. Tis is beacause the light beam spreads, as it propogates into a spherical region of area  $4\pi r^2$ , but LASER does not spread and hence its intensity remains constant.

**8.27** Electric field of an EM wave is an oscillating field and so is the electric force caused by it on a charged particle. This electric force averaged over an integral number of cycles is zero since its direction changes every half cycle. Hence, electric field is not responsible for radiation pressure.

**8.28** 
$$\mathbf{E} = \frac{\lambda \,\hat{\mathbf{e}}_s}{2\pi\varepsilon_o a} \,\hat{\mathbf{j}}$$

$$\mathbf{B} = \frac{\mu_o i}{2\pi a} \hat{\mathbf{i}}$$

$$=\frac{\mu_o \lambda v}{2\pi a}\hat{i}$$

$$\mathbf{S} = \frac{1}{\mu_o} (\mathbf{E} \times \mathbf{B}) = \frac{1}{\mu_o} \left( \frac{\lambda \, \hat{\mathbf{j}}_s}{2\pi\varepsilon_o a} \, \hat{\mathbf{j}} \times \frac{\mu_o \, \lambda v}{2\pi a} \, \hat{\mathbf{i}} \right)$$

$$=\frac{-\lambda^2 v}{4\pi^2 \varepsilon_0 a^2} \hat{\mathbf{k}}$$

8.29 Let the distance between the plates be d. Then the electric field  $E = \frac{V_o}{d} \sin(2\pi v t)$ . The conduction current density is given by the Ohm's law = E.

$$\Rightarrow J^{c} = \frac{1}{\rho} \frac{V_{o}}{d} \sin(2\pi vt) = \frac{V_{o}}{\rho d} \sin(2\pi vt)$$
$$= J_{o}^{c} \sin 2\pi vt$$

where 
$$J_0^c = \frac{V_0}{\rho d}$$
.

The displacement current density is given as

$$J^{d} = \varepsilon \frac{\partial E}{\partial t} = \varepsilon \frac{\partial}{\partial t} \left\{ \frac{V_{o}}{d} \sin(2\pi v t) \right\}$$
$$= \frac{\varepsilon 2\pi v V_{o}}{d} \cos(2\pi v t)$$

= 
$$J_0^d \cos(2\pi v t)$$
, where  $J_0^d = \frac{2\pi v \varepsilon V_0}{d}$ 

$$\int_{0}^{d} \int_{0}^{c} = \frac{2\pi v \varepsilon V_{o}}{d} \cdot \frac{\rho d}{V_{o}} = 2\pi v \varepsilon \rho = 2\pi \times 80 \varepsilon_{o} v \times 0.25 = 4\pi \varepsilon_{o} v \times 10$$

$$= \frac{10v}{9 \times 10^{9}} = \frac{4}{9}$$

**8.30** (i) Displacement curing density can be found from the relation be  $\mathbf{J}_D = \varepsilon_0 \frac{d\mathbf{E}}{dt}$ 

$$= \varepsilon_0 \,\mu_0 \,I_0 \,\frac{\partial}{\partial t} \cos (2\pi V t). \ln \left(\frac{s}{a}\right) \hat{\mathbf{k}}$$

$$= \frac{1}{c^2} I_0 2\pi v^2 \left(-\sin(2\pi v t)\right) \ln\left(\frac{s}{a}\right) \hat{\mathbf{k}}$$

$$= \left(\frac{v}{c}\right)^2 2\pi I_0 \sin(2\pi v t) \ln\left(\frac{a}{s}\right) \hat{k}$$

$$= \frac{2\pi}{\lambda^2} I_0 \ln \left(\frac{a}{s}\right) \sin(2\pi vt) \,\hat{\mathbf{k}}$$

(ii) 
$$I^d = \int J_D s ds d\theta$$

$$= \frac{2\pi}{\lambda^2} I_0 2\pi \int_{s=0}^a \ln\left(\frac{a}{s}\right) . s ds \sin\left(2\pi v t\right)$$

$$= \left(\frac{2\pi}{\lambda}\right)^2 I_0 \int_{s=0}^a \frac{1}{2} ds^2 \ln\left(\frac{a}{s}\right) \cdot \sin(2\pi vt)$$

$$= \frac{a^2}{4} \left(\frac{2\pi}{\lambda}\right)^2 I_0 \int_{s=0}^a d\left(\frac{s}{a}\right)^2 ln\left(\frac{a}{s}\right)^2 . \sin(2\pi vt)$$

$$= -\frac{\alpha^2}{4} \left(\frac{2\pi}{\lambda}\right)^2 I_0 \int_0^1 \ln \xi \, d\xi \cdot \sin(2\pi vt)$$

$$= + \left(\frac{a}{2}\right)^2 \left(\frac{2\pi}{\lambda}\right)^2 I_0 \sin 2\pi vt \quad (:: The integral has value -1)$$

$$I^{d} = \left(\frac{a}{2} \cdot \frac{2\pi}{\lambda}\right)^{2} I_{0} \sin 2\pi v t = I_{0}^{d} \sin 2\pi v t$$

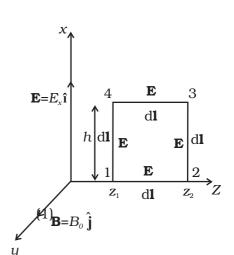
$$\frac{I_0^d}{I_0} = \left(\frac{a\pi}{\lambda}\right)^2.$$

**8.31** (i) 
$$\sqrt[8]{E}.dl = \int_{1}^{2} E.dl + \int_{2}^{3} E.dl + \int_{3}^{4} E.dl + \int_{4}^{1} E.dl$$

$$= \int_{1}^{2} E.dl \cos 90^{\circ} + \int_{2}^{3} E.dl \cos 0 + \int_{3}^{4} E.dl \cos 90^{\circ} + \int_{4}^{1} E.dl \cos 180^{\circ}$$

$$= \mathbf{E}_0 h[\sin(k\mathbf{z}_2 - \omega t) - \sin(k\mathbf{z}_1 - \omega t)]$$

(ii) For evaluating  $\int \mathbf{B.ds}$  let us consider the rectangle 1234 to be made of strips of area ds = h dz each.



 $\mathbf{B}=B_{u}\hat{\mathbf{j}}$ 

 $\int \mathbf{B.ds} = \int Bds \cos 0 = \int Bds = \int_{z_1}^{z_2} B_0 \sin(kz - \omega t) hdz$ 

$$= \frac{-B_o h}{k} \left[ \cos(kz_2 - \omega t) - \cos(kz_1 - \omega t) \right]$$
 (2)

(iii) 
$$\mathbf{\tilde{N}E.dl} = \frac{-d\phi_B}{dt}$$

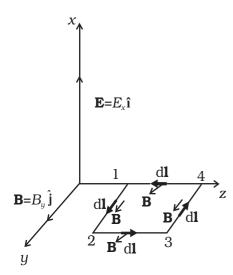
Using the relations obtained in Equations (1) and (2) and simplifying, we get

$$E_0 h[\sin(kz_2 - \omega t) - \sin(kz_1 - \omega t)] = \frac{B_o h}{k} \omega[\sin(kz_2 - \omega t) - \sin(kz_1 - \omega t)]$$

$$E_0 = B_0 \frac{\omega}{k}$$

$$\frac{E_0}{B_0} = c$$

(iv) For evaluating **\( \bar{\mathbb{B}}\).dl \( \), let us consider the loop 1234 in yz plane as shown in Fig.** 



$$\mathbf{\tilde{N}B.dl} = \int_{1}^{2} \mathbf{B.dl} + \int_{2}^{3} \mathbf{B.dl} + \int_{2}^{4} \mathbf{B.dl} + \int_{1}^{1} \mathbf{B.dl}$$

$$= \int_{1}^{2} B dl \cos 0 + \int_{2}^{3} B dl \cos 90^{\circ} + \int_{3}^{4} B dl \cos 180^{\circ} + \int_{4}^{1} B dl \cos 90^{\circ}$$

$$= B_0 h[\sin(kz_1 - \omega t) - \sin(kz_2 - \omega t)]$$
(3)

Now to evaluate  $\phi_E = \int {\bf E.ds}$ , let us consider the rectangle 1234 to be made of strips of area hdz each.

$$\phi_E = \int \mathbf{E}.\mathbf{ds} = \int Eds \cos 0 = \int Eds = \int_{Z_1}^{Z_2} E_0 \sin(kz_1 - \omega t) hdz$$

$$= \frac{-E_0 h}{k} \left[ \cos(kz_2 - \omega t) - \cos(kz_1 - \omega t) \right]$$

$$\therefore \frac{d\phi_E}{dt} = \frac{E_0 h \omega}{k} [\sin(kz_1 - \omega t) - \sin(kz_2 - \omega t)]$$

 $\mathbf{E} = E_x \hat{\mathbf{i}}$   $\mathbf{E} = B_y \hat{\mathbf{j}}$   $\mathbf{ds}$   $\mathbf{h}$   $\mathbf{dl}$ 

In 
$$\int_0^\infty \mathbf{B} \cdot d\mathbf{l} = \mu_0 \left( I + r_0 \frac{d\phi_E}{dt} \right)$$
,  $I = \text{conduction current}$   
= 0 in vacuum.

$$\therefore \mathbf{\tilde{N}} \mathbf{B}.\mathbf{dl} = \mu_0 \varepsilon_0 \frac{d\phi_E}{dt}$$

Using relations obtained in Equations (3) and (4) and ssimplifying, we get

(4)

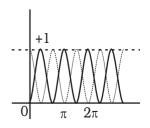
$$B_0 = E_0 \frac{\omega}{k} \mu \varepsilon_0$$

$$\frac{E_0}{B_0} \frac{\omega}{k} = \frac{1}{\mu \varepsilon_0}$$
 But  $E_0/B_0 = c$ , and  $\omega = ck$ 

or 
$$c.c = \frac{1}{\mu \varepsilon_0}$$
 Therefore,  $c = \frac{1}{\sqrt{\mu_0 \varepsilon_0}}$ .

**8.32** (a) E - field contribution is 
$$u_E = \frac{1}{2} \varepsilon_0 E^2$$

B - field contribution is 
$$u_B = \frac{1}{2} \frac{B^2}{\mu_0}$$



Total energy density 
$$u = u_E + u_B = \frac{1}{2}\varepsilon_0 E^2 + \frac{1}{2}\frac{B^2}{\mu_0}$$
 (1)

The values of  $E^2$  and  $B^2$  vary from point to point and from moment to moment. Hence, the effective values of  $E^2$  and  $B^2$  are their time averages.

$$(E^2)_{av} = E_0^2 [\sin^2(kz - \omega t)]_{av}$$

$$(B^2)_{av} = (B^2)_{av} = B_0^2 [\sin^2(kz - \omega t)]_{av}$$

The graph of  $\sin^2\theta$  and  $\cos^2\theta$  are identical in shape but shifted by  $\pi/2$ , so the average values of  $\sin^2\theta$  and  $\cos^2\theta$  are also equal over any integral multiple of  $\pi$ .

and also  $\sin^2\theta + \cos^2\theta = 1$ 

So by symmetry the average of  $\sin^2\theta$  = average of  $\cos^2\theta = \frac{1}{2}$ .

$$\therefore (E^2)_{av} = \frac{1}{2} E_0^2 \text{ and } (B^2)_{av} = \frac{1}{2} B_0^2$$

Substuting in Equation (1),

$$u = \frac{1}{4}\varepsilon_0 E^2 + \frac{1}{4}\frac{B_0^2}{\mu} \tag{2}$$

(b) We know 
$$\frac{E_0}{B_0} = c$$
 and  $c = \frac{1}{\sqrt{\mu} \varepsilon_0}$   $\therefore \frac{1}{4} \frac{B_0^2}{\mu_0} = \frac{E_0^2 / c^2}{4\mu_0} = \frac{E_0^2}{4\mu_0} \mu_0 \varepsilon_0 = \frac{1}{4} \varepsilon_0 E_0^2$ .

Therefore, 
$$u_{av} = \frac{1}{4}\varepsilon_0 E_0^2 + \frac{1}{4}\varepsilon_0 E_0^2 = \frac{1}{2}\varepsilon_0 E_0^2$$
, and  $I_{av} = u_{av}c = \frac{1}{2}\varepsilon_0 E_0^2$ .